

# ANIMATED PHASE PORTRAITS OF NONLINEAR AND CHAOTIC DYNAMICAL SYSTEMS

Jean-Marc Ginoux

► **To cite this version:**

Jean-Marc Ginoux. ANIMATED PHASE PORTRAITS OF NONLINEAR AND CHAOTIC DYNAMICAL SYSTEMS. A.H. Siddiqi, Pammy Manchada, M. Brokate. Calculus with Applications, I K International Publishing House Pvt. Ltd, pp.523-529, 2011, 9381141320. <hal-01056182>

**HAL Id: hal-01056182**

**<https://hal-univ-tln.archives-ouvertes.fr/hal-01056182>**

Submitted on 17 Aug 2014

**HAL** is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

# ANIMATED PHASE PORTRAITS OF NONLINEAR AND CHAOTIC DYNAMICAL SYSTEMS

JEAN-MARC GINOUX<sup>1</sup>

## 1. INTRODUCTION

The aim of this section is to present programs allowing to highlight the *slow-fast* evolution of the solutions of nonlinear and chaotic dynamical systems such as: Van der Pol, Chua and Lorenz models. These programs provide animated phase portraits in dimension two and three, i.e. “integration step by step” which are useful tools enabling to understand the dynamic of such systems.

## 2. VAN DER POL MODEL

The oscillator of B. Van der Pol [6] is a second-order system with non-linear frictions which can be written:

$$\frac{d^2y}{dt^2} - \alpha(1 - y^2)\frac{dy}{dt} + y = 0$$

The particular form of the friction which can be carried out by an electric circuit causes a decrease of the amplitude of the great oscillations and an increase of the small. This equation constitutes the “paradigm of relaxation-oscillations”. According to D’Alembert transformation [1] any single  $n^{th}$  order differential equation may be transformed into a system of  $n$  simultaneous first-order equations and conversely. Let’s consider that:  $\alpha(1 - y^2)y = \frac{d}{dt}\alpha(y - \frac{y^3}{3})$  and let’s pose:  $x_1 = y$  and  $y = -\alpha x_2$ . Thus, we have:

$$\begin{bmatrix} \dot{x}_1 \\ \dot{x}_2 \end{bmatrix} = \frac{1}{\alpha} \begin{bmatrix} x_1 + x_2 - \frac{x_1^3}{3} \\ -x_1 \end{bmatrix}$$

When  $\alpha$  becomes very large,  $x_1$  becomes a “fast” variable and  $x_2$  a “slow” variable. In order to analyze the limit  $\alpha \rightarrow \infty$ , we introduce a small parameter  $\varepsilon = 1/\alpha^2$  and a “slow time”  $t' = t/\alpha = \sqrt{\varepsilon}t$ . Thus, the system can be written:



All the main functions<sup>2</sup> are described in Tab. 1:

Property	Description
OutputFcn	A function for the solver to call after every successful integration step.
odephas2	2D phase plane
Refine	Increases the number of output points by a factor of Refine.

TABLE 1. Description of the main functions

The program provides the animated phase portrait corresponding to the solution of system (1) and the time series.

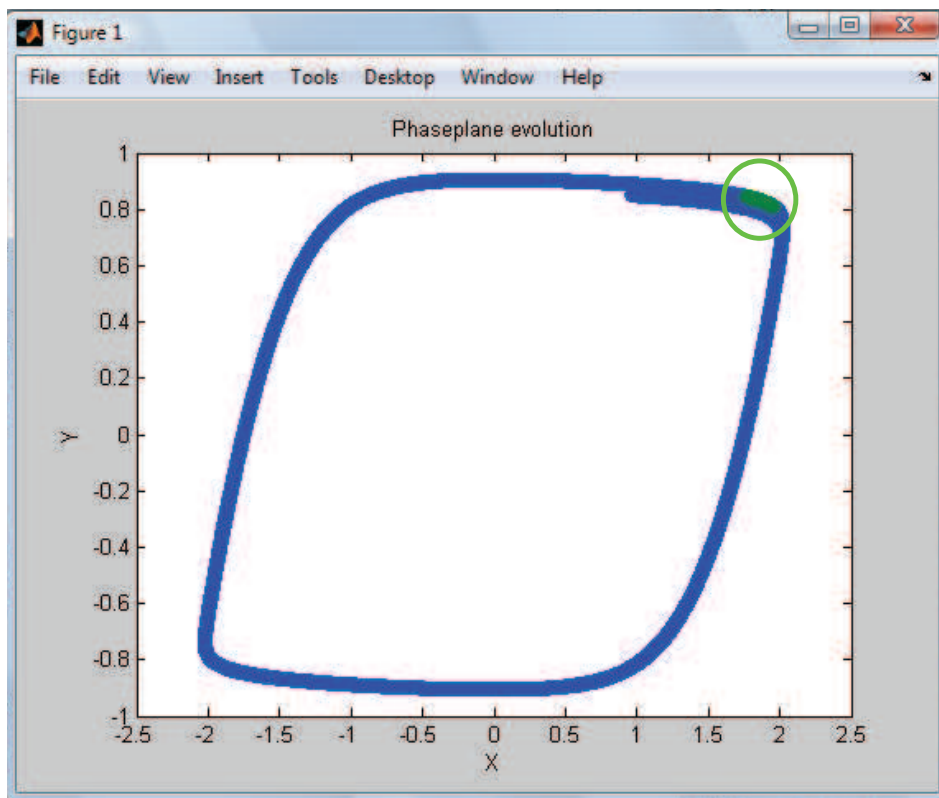


FIGURE 2. Animated phase portrait in 2D

<sup>2</sup>See the textbook of Prof. Rudra Pratap [5]

### 3. CHUA'S MODEL

The L.O. Chua's circuit [2] is a relaxation oscillator with a cubic non-linear characteristic elaborated from a circuit comprising a harmonic oscillator of which operation is based on a field-effect transistor, coupled to a relaxation-oscillator composed of a tunnel diode. The modeling of the circuit uses a capacity which will prevent from abrupt voltage drops and will make it possible to describe the fast motion of this oscillator by the following equations which also constitute a *slow-fast dynamical system* or a *singularly perturbed dynamical system*.

$$(2) \quad \begin{bmatrix} \varepsilon \dot{x}_1 \\ \dot{x}_2 \\ \dot{x}_3 \end{bmatrix} = \begin{bmatrix} x_3 - \frac{44}{3}x_1^3 - \frac{41}{2}x_1^2 - \mu x_1 \\ -x_3 \\ -0.7x_1 + x_2 + 0.24x_3 \end{bmatrix}$$

with  $\varepsilon$  and  $\mu$  are real parameters  $\varepsilon = 0.05$ ,  $\mu = 2$ . The system (2) which can not be integrated by quadratures (closed-form) exhibits a solution evolving on "chaotic attractor" in the shape of a "double-scroll". The program presented here enables to emphasize the *slow-fast* evolutions of the solution on the "chaotic attractor".

First, copy the files named "chua" and "chua1" into the "current folder" of MatLab. Then, open the m-file called "chua1" and press the green button to provide an animated plot 3D. The solution materialized by a green point evolves on the attractor according to slow and fast motion. The function "odephas2" is simply replaced by "odephas3".

### 4. LORENZ MODEL

The purpose of the model established by Edward Lorenz [4] was in the beginning to analyze the unpredictable behavior of weather. After having developed non-linear partial derivative equations starting from the thermal equation and Navier-Stokes equations, Lorenz truncated them to retain only three modes. The most widespread form of the Lorenz model is as follows:

$$(3) \quad \begin{bmatrix} \dot{x}_1 \\ \dot{x}_2 \\ \dot{x}_3 \end{bmatrix} = \begin{bmatrix} \sigma(x_2 - x_1) \\ -x_1x_3 + rx_1 - x_2 \\ x_1x_2 - \beta x_3 \end{bmatrix}$$

with  $\sigma$ ,  $r$ , and  $\beta$  are real parameters:  $\sigma = 10$ ,  $\beta = \frac{8}{3}$ ,  $r = 28$ .

Although, this system is not *singularly perturbed* since it does not contain any small multiplicative parameter, it is a *slow-fast dynamical system*. Its solution exhibits a solution evolving on “chaotic attractor” in the shape of a “butterfly”. The program presented here enables to emphasize the *slow-fast* evolutions of the solution on the “chaotic attractor”.

Name	Description
vanderpol	Dynamical system (1) with $\varepsilon = 1/20$ .
vanderpolpp	animated phase portrait 2D
chua	Dynamical system (2) with $\varepsilon = 1/20$ , $\mu = 2$ .
chua1	animated phase portrait 3D
lorenz	Dynamical system (3) with $\sigma = 10$ , $\beta = 8/3$ $r = 28$
lorenz1	animated phase portrait 3D

TABLE 2. List of program for animated phase portraits.

#### REFERENCES

- [1] J. D’ALEMBERT, Suite des recherches sur le calcul intégral, quatrième partie : Méthodes pour intégrer quelques équations différentielles, *Hist. Acad. Berlin*, tome IV (1748) 275-291.
- [2] L.O. CHUA, M. KOMURO & T. MATSUMOTO, “The double scroll family,” *IEEE Trans. on Circuits and Systems*, 33(11) (1986) 1072-1097.
- [3] J.M. GINOUX, *Differential geometry applied to dynamical systems*, World Scientific Series on Nonlinear Science, Series A **66** (World Scientific, Singapore), 2009.
- [4] E. N. LORENZ, “Deterministic non-periodic flows,” *Journal of Atmospheric Sciences*, 20 (1963) 130-141.
- [5] R. PRATAP, *Getting Started with MATLAB: A Quick Introduction for Scientists and Engineers*, Oxford University Press, USA (November 16, 2009).
- [6] B. VAN DER POL, “On ‘Relaxation-Oscillations’,” *Philosophical Magazine*, 7(2) (1926) 978-992.

<sup>1</sup> LABORATOIRE PROTEE, I.U.T. DE TOULON, SENIOR LECTURER IN APPLIED MATHEMATICS, DOCTOR IN MATHEMATICS, DOCTOR IN HISTORY OF SCIENCES, UNIVERSITÉ DU SUD, BP 20132, F-83957 LA GARDE CEDEX, FRANCE  
*E-mail address:* ginoux@univ-tln.fr, <http://ginoux.univ-tln.fr>